

## Manuscript Details

<b>Manuscript number</b>	YMEHY_2018_150
<b>Title</b>	A proposed tandem mechanism for memory storage in neurons involving magnetite and prions
<b>Article type</b>	Full length article

### Abstract

Knowledge about how information is stored in neurons of animals and in the human brain is still incomplete. A hypothesis related to long-term changes in synaptic efficiency has strong experimental support, but does not seem to be able to explain all observations. It has recently been proposed that magnetite together with a prion-like protein could be involved in a tandem mechanism for storage of memory in neurons in which electric impulses are received and reshaped by the magnetite to a form which can be accepted by the protein. The magnetite crystals can be magnetized by an electrical impulse, but it cannot hold the magnetism, which drops to zero after each impulse. Therefore, magnetite cannot be the substance in which information is stored. In the present paper we explain how a tandem mechanism could function in a neuron in which magnetite is situated together with a prion-like protein close to the cell membrane of the axon. We assume in addition that the information is stored in special storage neurons. With this we propose a new hypothesis for information storage in neurons which could operate in addition to synaptic plasticity, but perhaps in different neurons.

<b>Taxonomy</b>	Natural Sciences, Health Sciences
<b>Corresponding Author</b>	Lars Walloe
<b>Order of Authors</b>	Erik Alfsen, Fredrik Størmer, Arild Njå, Lars Walloe
<b>Suggested reviewers</b>	Erica White-Grindley, Michele Pignatelli, Kausik Si

## Submission Files Included in this PDF

### File Name [File Type]

Tandem mechanism Letter to Editor MedHyp 07.02.18.docx [Cover Letter]

Tandem mechanism MedHyp 070218 LW.docx [Manuscript File]

Conflict of interest.docx [Conflict of Interest]

Alfsen.pdf [Supporting File]

To view all the submission files, including those not included in the PDF, click on the manuscript title on your EVISE Homepage, then click 'Download zip file'.



**UNIVERSITY OF OSLO**  
FACULTY OF MEDICINE

**Lars Walløe**

Professor emeritus  
Department of Physiology  
Institute of Basic Medical Sciences  
P.O.Box 1103 Blindern  
0317 Oslo  
Norway

Sognsvannsveien 9

Oslo, 7 February 2018

Phone: +47 22 85 12 18

Fax: +47 22 85 12 49

E-mail: [lars.walloe@medisin.uio.no](mailto:lars.walloe@medisin.uio.no)

<http://folk.uio.no/larswa/>

Editor,  
Medical Hypotheses

Manuscript: **A proposed tandem mechanism for memory storage in neurons involving magnetite and prions**

In our submitted manuscript, we propose a tandem mechanism for memory storage which involves magnetite and prions. The magnetization of magnetite can only be explained by a quantum mechanical process. The quantum mechanical argument is probably too detailed to be included in the article. However, if our manuscript is accepted for publication in Medical Hypotheses, we shall archive a six page note on the magnetization of magnetite chains in neurons in Cornell University Library: **arXiv.org** - in its section for Mathematical Physics. The note is written by our first author Erik M. Alfsen. I have attached this note on the magnetization of magnetite chains in neurons to this letter in case you would like to attach it as supplementary information to potential reviewers of the submitted manuscript.

On behalf of the authors,

Yours sincerely

Lars Walløe

1 *Manuscript title:*

2 A proposed tandem mechanism for memory storage in neurons involving  
3 magnetite and prions

4

5 *Abbreviated title:*

6 Memory storage with magnetite and prions

7

8 Erik M. Alfsen, Department of Mathematics, University of Oslo, Norway

9 Fredrik C. Størmer, Norwegian Institute of Public Health, Oslo, Norway

10 Arild Njå, Division of Physiology, Institute of Basic Medical Sciences, University  
11 of Oslo, Norway

12 Lars Walløe, Division of Physiology, Institute of Basic Medical Sciences,  
13 University of Oslo, Norway

14

15 Correspondence should be addressed to LW (lars.walloe@medisin.uio.no)

16

17 Submitting and Corresponding author:

18 Professor Lars Walløe MD PhD

19 Division of Physiology,

20 Institute of Basic Medical Sciences

21 University of Oslo

22 P.O.Box 1103 Blindern

23 NO-0317 Oslo

24 Norway

25

26 Authors report no conflict of interest

27

28 The research is funded by the two university departments

29

30 **Abstract**

31

32 Knowledge about how information is stored in neurons of animals and  
33 in the human brain is still incomplete. A hypothesis related to long-term  
34 changes in synaptic efficiency has strong experimental support, but does not  
35 seem to be able to explain all observations. It has recently been proposed that  
36 magnetite together with a prion-like protein could be involved in a tandem  
37 mechanism for storage of memory in neurons in which electric impulses are  
38 received and reshaped by the magnetite to a form which can be accepted by  
39 the protein. The magnetite crystals can be magnetized by an electrical  
40 impulse, but it cannot hold the magnetism, which drops to zero after each  
41 impulse. Therefore, magnetite cannot be the substance in which information  
42 is stored. In the present paper we explain how a tandem mechanism could  
43 function in a neuron in which magnetite is situated together with a prion-like  
44 protein close to the cell membrane of the axon. We assume in addition that  
45 the information is stored in special *storage neurons*. With this we propose a  
46 new hypothesis for information storage in neurons which could operate in  
47 addition to synaptic plasticity, but perhaps in different neurons.

48

49 **Keywords:** memory storage, magnetite, prions

50

51 **Introduction**

52

53 Knowledge of the mechanisms for storage of memory in neurons is  
54 incomplete. A hypothesis related to long-term changes in synaptic strength  
55 and the growth of new synaptic connections has strong experimental support  
56 (Kandel, 2001). In addition, nonsynaptic processes, such as changes in  
57 neuronal excitability, may contribute to some forms of learning and memory  
58 (Mozzachiodi and Byrne, 2010). The formation of long-term memories  
59 requires protein synthesis, and the gene regulation involved in such processes  
60 has been shown to depend on special protein kinases (Giese and Mizuno,  
61 2013) and more recently, on prion-like proteins (Si and Kandel, 2016) .

62 Prion-like proteins (here referred to as “prions” for short) have the  
63 crucial property that they can change shape from a normal resting form to an  
64 alternate chain-formed version. For brevity, we refer to a prion in the  
65 normal form as being in the “ground state”, and to one in the alternate form  
66 as being in the “excited state”. In Si and Kandel (2016) it is explained how a  
67 prion can be switched from the ground state to the excited state by an  
68 intricate *prion-chain reaction* triggered by an electric impulse.

69 It has also been proposed that magnetite is involved in the storage of  
70 information in neurons (Størmer and Laane, 2009, and Banaclocha et al,  
71 2010). Magnetite is an iron oxide ( $\text{Fe}_3\text{O}_4$ ) which is widely found in living  
72 organisms without being involved in any known biochemical reactions.

73 Electron microscopy has indicated the presence of a minimum of 5  
74 million single-domain magnetite crystals per gram in human brain tissue.  
75 Black strings of aggregated particles can be extracted from brain tissue and  
76 can be viewed under low power through an optical dissecting microscope. The  
77 crystals range in size between 10 and 70 nm in diameter (Kirschvink et al,  
78 1992). Magnetite crystals have also been found in bacteria and fish, where  
79 they were seen to be organized in membrane-bound chains with up to 80  
80 single-domain crystals per chain (Kirschvink et al, 1985, 2001). Such chains  
81 appear to be present across all animal species.

82 The particles in these chains are formed of magnetite crystals that are  
83 different from those formed through geological processes. The crystal  
84 morphology of magnetite in living organisms, referred to as *nanocrystalline*  
85 *magnetite*, is in fact particularly well suited for information storage for various  
86 reasons. Here, we only wish to point out that nanocrystalline magnetite has  
87 the crucial property that it can easily be magnetized by electrical impulses, but  
88 cannot hold the magnetism, which drops to zero after each impulse.

89 Therefore, nanocrystalline magnetite cannot be the substance in which  
90 information is stored. But it has recently been suggested that magnetite  
91 chains together with prions could be involved in a tandem mechanism in  
92 which incident electric impulses are received and reshaped by the magnetite,  
93 giving it a form that can be accepted by prion-like proteins in which the  
94 information can be permanently stored (Størmer and Alfsen, 2015).

95 In this paper we will explain how such a mechanism may function in a  
96 neuron where a nanocrystalline magnetite chain is located together with a  
97 prion close to the cell membrane, and how the electric impulses could trigger  
98 prion-chain reactions such that the information is stored in the prion.

99

100

### 101 **A storage cell hypothesis**

102

103 Nerve impulses (action potentials) are all-or-none electrical signals  
104 that encode neural information for long-distance communication along  
105 nerve fibers (axons). Action potentials have large, fixed amplitudes (about  
106 100 mV), short durations (about 1 ms) and travel through long axons at  
107 speeds that range from below 1 to above 100 m s<sup>-1</sup>, depending on the axon  
108 diameter and degree of myelination.

109           Action potentials are elicited if the membrane potential is depolarized  
110 to a threshold level. The threshold is lowest in the initial segment of the axon.  
111 After an action potential is generated at this site, it can spread without decay  
112 down the axon and all of its branches, and may (although with signal decay)  
113 also invade the cell body and proximal dendrites. Graded transmembrane  
114 potentials show exponential decay with distance from the generation site in  
115 the neuron, but multiple small potentials may summate so that the threshold  
116 is reached and an action potential is fired in the initial segment of the axon,  
117 even if each input is too small to reach the threshold alone.

118           The transmembrane ionic currents associated with the passage of a  
119 propagated action potential (an inward  $\text{Na}^+$  current through voltage-gated  
120 sodium channels, and, after a brief delay, an outward  $\text{K}^+$  current through  
121 voltage-gated potassium channels) will generate a local rotating current that  
122 could potentially magnetize a properly oriented rod of magnetite  
123 nanocrystals, as we explain here. We propose a tandem mechanism by which  
124 the propagation of an action potential in an axon could, given that specific  
125 conditions are satisfied, magnetize a local rod of magnetite nanocrystals that  
126 in turn could change the conformation of a local prion and store a permanent  
127 memory of the event.

128



129 *How can electrical signals magnetize the magnetite chain?*

130       According to the tandem hypothesis, information is stored in neurons  
131 (here referred to as *storage neurons*) which contain a magnetite chain  
132 together with a prion that can be in one of two possible states. More  
133 specifically, each storage neuron is assumed to contain just one magnetite  
134 chain, and one bit of information is stored in the neuron when the prion is  
135 switched from the ground state to the excited state.

136       In a storage neuron, the bursts of  $\text{Na}^+$  ions into the cell and  $\text{K}^+$  ions out  
137 of the cell create vortices that rotate on their axes like short-lived electrical  
138 currents. When such vortices rotate around a magnetite chain, they can  
139 magnetize it. These vortices arise when two oppositely directed flows of ions  
140 pass through neighboring membrane areas. If P1 and P2 are sufficiently close  
141 ion channels, the two flows will interact and produce vortices that rotate  
142 clockwise in the (x-z)-plane (Figure 1).

143       According to this argument, incident signals in a storage neuron can  
144 create vortices of short-lasting electrical currents that rotate like a whirlwind  
145 around the magnetite chain, and these currents can magnetize the chain.  
146 However, the magnetite chain differs from an electromagnet, as described  
147 below.

148

149 *How can the magnetite chain change the shape of the prion?*

150           We have seen that the magnetite chain in a storage neuron can be  
151 magnetized by currents rotating around it, but it is not so easy to explain how  
152 the magnetic chain can change the shape of the prion, since it requires an  
153 electric and not a magnetic impulse to do this. Here it may be helpful to use  
154 an analogy with a familiar electric device, a transformer, which sets up a  
155 pulsating electromagnetic field that produces eddy currents.

156           In the same way, when the magnetite chain of a storage neuron is  
157 magnetized by electric pulses, it will set up a pulsating electric field which  
158 produces ionic currents. These will rotate in right-handed vortices around the  
159 magnetized magnetite chain like a whirlwind. A detailed proof of this can be  
160 obtained by an elementary argument, but all we need to know here is that the  
161 electric field at a given point is largest close to the magnetized chain, and that  
162 it decreases rapidly with increasing distance from the chain.

163           It is thus conceivable that a magnetite chain can indeed act like a pulse  
164 transformer, and “assemble” and “coordinate” many small electric currents  
165 around a magnetite chain which is located together with a prion close to the  
166 cell membrane, and that it can replace these currents by a strong current  
167 which rotates in the same way around the magnetite chain and is strong  
168 enough to magnetize it.

169 But this needs a more detailed explanation. For this, we need to know  
170 more about the magnetic properties of the magnetite chains, which are quite  
171 different from the magnetic properties of iron cores in standard transformers.  
172

173 *Why is nanocrystalline magnetite particularly suited for the storage of*  
174 *information in neurons?*

175 Magnetite ( $\text{Fe}_3\text{O}_4$ ) is an iron oxide which differs from ordinary rust ( $\text{Fe}_2$   
176  $\text{O}_3$ ) in its crystalline structure and its specific magnetic properties. In  
177 molecules of ordinary rust, all the innermost electronic shells are filled, but in  
178 molecules of magnetite, there is one loosely bound electron outside the full  
179 innermost shells. The extraordinary magnetic properties of magnetite are  
180 largely due to this electron. The crystal structure of nanocrystalline magnetite  
181 in neurons differs from that of magnetite in geological deposits. Geological  
182 magnetite crystals are usually octahedral, whereas those of biological  
183 magnetite are cubo-octahedral. The individual crystals in a magnetite chain  
184 are tiny magnets which are aligned “north pole to south pole” along the axis  
185 of the chain, each with its poles in the faces of the crystal lying perpendicular  
186 to the axis.

187 We know that magnetization of ferromagnetic substances involves  
188 crystal *domains*, each usually containing a large number of crystals, and that

189 the magnetization process consists of a rotational displacement of these  
190 domains so that they are oriented in the direction of the magnetizing  
191 magnetic field (Harnwell (1949)).

192       The magnetite chain in a storage neuron cannot be magnetized by  
193 rotation of its single crystal domains since they are tiny magnets and are  
194 already oriented north pole to south pole along the chain axis. However, they  
195 can be magnetized by a quantum mechanical process. This is based on  
196 quantum mechanical resonance between a strong magnetizing exterior field  
197 and the weak magnetic fields of the crystals in the chain, which are forced to  
198 oscillate in the same quantum mechanical phase as the magnetizing field. The  
199 process is explained in more detail in Alfsen (2018) where the phases in  
200 question are related to the complex wave functions of quantum mechanics.  
201 Here we explain the idea in non-mathematical terms. For this purpose we will  
202 make use of the old quantum theory that the electrons in an atom rotate in  
203 orbits around the atomic nucleus. In this picture, the loosely bound electron in  
204 the magnetite molecule may be thought of as a particle rotating in orbit  
205 around a 'core' consisting of the atomic nuclei of iron and oxygen and the  
206 other tightly bound electrons.

207       When a storage neuron is in the resting state, the crystals in its  
208 magnetite chain are completely unrelated in phase. But when the chain is

209 subject to a sufficiently strong exterior magnetic field, all crystals in the chain  
210 are forced to oscillate in the same phase as the exterior field. After this, the  
211 chain can no longer contain crystals in opposite phase whose magnetic fields  
212 erase each other out. Therefore, the crystals are now all in the same phase as  
213 the magnetizing exterior field, and the chain is magnetized up to its greatest  
214 possible strength.

215         The magnetization of a magnetite chain by this process requires a  
216 certain critical energy level. If the critical energy level is exceeded, the chain  
217 will always be magnetized up to the same maximal strength. The critical  
218 energy level is here called the *magnetization energy* of the chain. Thus, no  
219 magnetization of the magnetite chain of a storage neuron will occur if it is  
220 subjected to energy levels below the magnetization energy, whereas any  
221 energy level above this magnetization energy will magnetize it up to its full  
222 strength but no further, however strong the signal is. This means that the  
223 magnetization process of the magnetite chain in a storage neuron is different  
224 from that of the iron core in an electromagnet, in which the magnetism  
225 increases proportionally with the strength of the magnetizing current.

226         Each domain in the nanocrystalline magnetite in neurons contains a  
227 single crystal. These crystals are of the same shape and size, and they also act  
228 as tiny magnets. When the magnetic chain of a storage neuron is magnetized

229 up to its maximal strength, the magnetic energy of the chain is equal to the  
230 magnetic energy of a single crystal multiplied by the number of crystals in the  
231 chain. The magnetization energy of a magnetite chain therefore depends only  
232 on its length, i.e. on the number of crystals in the chain.

233         According to the tandem hypothesis, one bit of information is stored  
234 when the prion in a storage neuron is switched from the ground state to the  
235 excited state. This requires a certain amount of energy which must be  
236 provided by the magnetized magnetite chain. We call this energy the  
237 *excitation energy* of the prion.

238         In order to raise the prion from the ground state to the excited state,  
239 the incident signal must provide sufficient energy to magnetize the chain, and  
240 the chain must in turn provide enough energy to create “electrical  
241 whirlwinds” that can excite the prion. This means that the magnetization  
242 energy held by the magnetite chain, which depends on its length, must exceed  
243 the excitation energy of the prion. The necessary amount of energy can  
244 therefore only be provided by a magnetite chain of a certain length. The  
245 minimum chain length needed for excitation of the prion is here called the  
246 *critical chain length* for the magnetic chain in a given storage neuron.

247         The discussion above shows that there are two conditions which must  
248 be satisfied by a storage neuron in a pathway. The first pertains to the storage

249 neuron itself, which must be able to provide the magnetization energy needed  
250 to magnetize the magnetite chain to its maximum strength, and the second  
251 pertains to the magnetite chain, which must contain enough crystals to be  
252 above the critical chain length.

253         The quantum mechanical process for magnetization of nanocrystalline  
254 magnetite requires much less energy than the magnetization of geological  
255 magnetite, which involves physical displacement of large domains. This  
256 explains why nanocrystalline magnetite has the crucial property that it is easy  
257 to magnetize but cannot hold the magnetism.

258

259

260 **Discussion**

261

262         How can the proposed tandem mechanism function in neurons?

263         White-Grindley et al. (2014) point out that: “- - - a central question is  
264 how the conversion of neuronal CPEB /cytoplasmic polyadenylation element  
265 binding protein/ to the amyloidogenic state is regulated to confer activity  
266 dependence and restrict it to the relevant neuron/ synapse.” In our tandem  
267 hypothesis context, this is the question of how the process that changes the

268 state of a prion is regulated to take place at just the right time and the right  
269 place in the relevant neuron or synapse.

270 It has been an open question whether memory storage takes place only  
271 in synapses or also in specialized neurons. We assume that both entities may  
272 be involved in the storage process. If information storage is based on the  
273 summation of a large number of local signals, the important entities are the  
274 synapses at which the incident signals are targeted. But this is not the case for  
275 information storage in storage neurons based on the proposed tandem  
276 mechanism. In this case, the most important entities are not the synapses but  
277 the initial segments of the axons of the post-synaptic neurons where action  
278 potentials are initiated and where magnetite chains are presumed to be  
279 located.

280 Here we discuss the propagating action potentials in a pathway of axons  
281 which terminates in a storage neuron. The magnetite chain in a storage  
282 neuron can be magnetized most efficiently if it is located laterally and near  
283 the cell membrane. We say that the storage neuron has a “threshold density  
284 of ion channels” if there are enough  $\text{Na}^+$  and  $\text{K}^+$  channels near the magnetite  
285 chain to release the rotating currents of ions that can magnetize the chain.

286 We argue, based on the quantum mechanical properties of  
287 nanocrystalline magnetite, that the *magnetization energy* of the magnetite



288 chain depends only on the length of the chain, which must be above the  
289 *critical length* to provide the *excitation energy* needed to switch the prion in  
290 the cell from the ground state to the excited state.

291 The following three assumptions are required as part of the tandem  
292 hypothesis for memory storage by magnetite and prions:

293 A1. The magnetite chain in a storage neuron is located close to the cell  
294 membrane and is laterally oriented in the neuron.

295 A2. The density of ion channels in the storage neuron is above the  
296 threshold level.

297 A3. The length of the magnetite chain in the storage neuron is above  
298 the critical length.

299 Given these assumptions, the tandem hypothesis explains how a  
300 propagating signal can magnetize the magnetite chain in a storage neuron,  
301 which in turn raises the prion in the storage neuron to the excited state and  
302 thereby stores one bit of information.

303 The proposed tandem mechanism converts electrical signals in the form  
304 of action potentials and propagated along axon pathways to a form that can  
305 be accepted by special “storage neurons”, each containing a chain of  
306 nanocrystalline magnetite and a prion which can be switched from a “ground

307 state” to an “excited (amyloidogenic) state” when the chain is magnetized by  
308 incident signals. This process stores one bit of information in the cell.

309

310 *Conclusions*

311       According to the tandem hypothesis, the storage neurons must satisfy  
312 three assumptions. The first of these specifies the location of the magnetite  
313 chain in the cell, the second requires a sufficient number of Na<sup>+</sup> and K<sup>+</sup> ion  
314 channels near the magnetite chain to permit the ion bursts that magnetize the  
315 chain, and the third requires the magnetite chain to be long enough to  
316 provide sufficient energy to raise the prion to the excited state and store  
317 memory in the cell

318       If cells are to function as storage neurons, it must also be possible to  
319 retrieve stored information when needed, through processes that are  
320 dependent on whether the prion is in the excited state or the ground state. It  
321 is conceivable that such retrieval of information could be achieved by sending  
322 a test signal into the cell via an appropriate synaptic connection.

323

324

325 **References**

326

327 Alfsen EM (2018) A note on magnetization of magnetite chains in neurons.

328 To be archived in Cornell University Library: arXiv.org – section for

329 Mathematical Physics

330

331 Banaclocha MAM, Bókkon I, Banaclocha HM (2010) Long term memory in

332 brain magnetite. Medical Hypotheses 74:254-7.

333

334 Giese KP, Mizuno, K (2013) The roles of protein kinases in learning and

335 memory. Learn Mem. 20:540-52.

336

337 Harnwell GP (1949) Principles of Electricity and Magnetism. Chap.11,

338 Sec.11.6, McGraw-Hill.

339

340 Kandel ER (2001) The molecular biology of memory storage: A dialogue

341 between genes and synapses. Science 294:1030-38.

342

- 343 Kirschvink JL, Walker MM, Chang S-B, Dizon AE, et al. (1985) Chains of  
344 single-domain magnetite particles in chinook salmon, *Oncorhynchus*  
345 *tshawytscha*. J. of Comparative Physiol. 157: 375-81.  
346
- 347 Kirschvink JL, Kobayashi-Kirschvink A, Woodford BJ (1992) Magnetite  
348 biomineralization in the human brain. Proc. Natl. Acad. Sci. USA 89: 7683-87.  
349
- 350 Kirschvink JL, Walker MM, Diebel CE (2001) Magnetite-based  
351 magnetoreception. Current opinion in Neurobiology 11; 462-67.  
352
- 353 Mozzachiodi R, Byrne JH (2010) More than synaptic plasticity: Role of  
354 nonsynaptic plasticity in learning and memory. Trends Neurosci 33:17-26.  
355
- 356 Si K, Kandel, ER (2016) The role of functional prion-like proteins in the  
357 persistence of memory. Cold Spring Harb Perspect Biol. 8:a021774.  
358
- 359 Størmer FC, Alfsen EM (2015) Is a combination of magnetite and prions  
360 involved in the memory storage in the human brain? Medical Hypotheses 85:  
361 111.  
362

- 363 Størmer FC, Laane CMM (2009) Membrane-bound memory in the neurons?  
364 Medical Hypotheses 73: 462.
- 365
- 366 White-Grindley E, Li L, Khan RM, Ren F, et al. (2014) Contribution of Orb2a  
367 stability in regulated amyloid-like oligomerization of *Drosophila* Orb2. PLOS  
368 Biol Vol.12| 12:e1001786.
- 369
- 370
- 371

372

373 Figure 1

374

375 Two oppositely directed bursts of ions through neighboring ion channels can  
376 create vortices that magnetize the magnetite chain in a storage neuron.

377 The action potential travels along the cell membrane in the x-direction.

378 Near point  $P_2$  the sodium channels are open and near the point  $P_1$  the

379 potassium channels. The magnetite chain S-N is posed laterally close to the

380 cell membrane and oriented in the y-direction.

381



The authors report no conflict of interest



# A note on magnetization of magnetite chains in neurons

Erik M. Alfsen

February 3, 2018

In article [1], it is proposed that information can be stored in a neuron when a magnetized chain of magnetite crystals is reshaped and amplified by electric impulses that are switching a prion from its ground state to its excited state. In this article, this process is explained without use of mathematics. But in the present note, we will present the quantum mechanical equations for the process in which such a chain of crystals is magnetized for this purpose.

We know that quantal quantities are denoted by complex-valued functions  $\psi = re^{i\varphi}$  defined on 4-dimensional space-time (or more general real spaces not needed here), and that the absolute value  $|\psi|$  gives the probability density of the result of measurements of such quantities. We also know that quantum mechanical superposition of two quantal quantities is represented by the sum of their complex wave functions. Actually, each quantum mechanical wave function  $\psi$  consists of two spin-components, here denoted by  $\psi_+$  and  $\psi_-$ . The mathematical theory of spin dates back to Pauli's 1927-article: On the quantum mechanics of the magnetic electron [2]. Here we will give a brief survey, based on the presentation in the book [3], of those parts of Pauli's theory on electron spin which will be needed in the present note.

From the two spin-components  $\psi_+$  and  $\psi_-$  we can construct a two-dimensional vector  $\Psi$  (not in physical space but in two-dimensional complex space) and the vector conjugate to it

$$\Psi = \begin{pmatrix} \psi_+ \\ \psi_- \end{pmatrix}, \quad \Psi^\dagger = (\psi_+^*, \psi_-^*). \quad (1)$$

As explained in [2; Sec. 10.1], the magnetic moment-vector of Pauli's magnetic electron is given by an equation of the form

$$\boldsymbol{\mu} = \mu \boldsymbol{\Sigma}, \quad (2)$$

where  $\mu$  is a constant (the scalar magnetic moment), and  $\Sigma$  is a three-dimensional vector (in physical space) given by the following equation for the coordinates

$$\Sigma_i = \Psi^\dagger \sigma_i \Psi, \quad i = 1, 2, 3 \quad \text{or} \quad i = x, y, z, \quad (3)$$

where  $\sigma_i$  are the *Pauli matrices*

$$\sigma_1 = \sigma_x = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \sigma_2 = \sigma_y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \quad \sigma_3 = \sigma_z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}. \quad (4)$$

(Here (1) is eq. (10.16) and (3) is eq. (10.12) in [3].)

Multiplying out the matrix product in (3) gives the following equations (which are (10.6a), (10.6b), (10.6c) in [3])

$$\Sigma_x = \psi_+^* \psi_- + \psi_-^* \psi_+ \quad (5a)$$

$$\Sigma_y = -i(\psi_+^* \psi_- - \psi_-^* \psi_+) \quad (5b)$$

$$\Sigma_z = \psi_+^* \psi_+ - \psi_-^* \psi_- . \quad (5c)$$

The spin-components of Pauli's magnetic electron are supposed to transform as follows under rotation by angle  $\alpha$  around a straight line in space

$${}'\psi_+ = e^{-i\alpha/2} \psi_+, \quad {}'\psi_- = e^{i\alpha/2} \psi_- . \quad (6)$$

By these equations, the wave function changes sign after rotation by the angle  $\alpha = 2\pi$ . But this won't cause problems, as it is not the wave function itself that has physical meaning, but quantities constructed from it, which are invariant under a sign-change.

The key example is the squared norm of the associated two-dimensional vector

$$|\Psi|^2 = (\psi_+^*, \psi_-^*) \begin{pmatrix} \psi_+ \\ \psi_- \end{pmatrix} = |\psi_+|^2 + |\psi_-|^2, \quad (7)$$

which is interpreted as the probability density of finding the electron with given spin-state at a given point in space-time [3; eq. (10.5)]. What we are here interested in, is the transformation of the magnetic moment-vector of an electron whose wave-function is transformed as in (6).

We know that the components of the vector  $\Sigma$ , which determines the magnetic moment-vector of the electron, is given as in (5). From this, we first observe that the  $z$ -component of  $\Sigma$  is transformed as follows

$${}'\Sigma_z = \psi_+^* \psi_+ - \psi_-^* \psi_- = \Sigma_z . \quad (8)$$

Thus, the  $z$ -component of  $\Sigma$  remains unchanged under this transformation of the wave function. This fact is of course no mere coincidence: the transformation formula (6) was tailored to achieve just this!

To determine the transformation of  $\Sigma_x$  and  $\Sigma_y$ , we assume that the wave function of the electron at a given point  $(x, y, z)$  at a given time, say  $t = 0$ , is given by the equation

$$\psi_+ = r_+ e^{i\varphi_+}, \quad \psi_- = r_- e^{i\varphi_-}, \quad (9)$$

with the *amplitudes*  $r_+$ ,  $r_-$  and the *phase angles*  $\varphi_+$  and  $\varphi_-$ .

Note that the two terms  $\psi_+^* \psi_-$  and  $\psi_-^* \psi_+$  occurring in (5a) and (5b) are given as follows in terms of the *phase difference*  $\varphi = \varphi_+ - \varphi_-$

$$\psi_+^* \psi_- = r_+ r_- e^{-i\varphi}, \quad \psi_-^* \psi_+ = r_+ r_- e^{i\varphi},$$

and that the two terms  $\psi_+^* \psi_+$  and  $\psi_-^* \psi_-$  occurring in (5c) are given by

$$\psi_+^* \psi_+ = r_+^2, \quad \psi_-^* \psi_- = r_-^2.$$

By (6) the spin-components of an electron transform as follows under rotation by the angle  $\alpha$

$$' \psi_+ = r_+ e^{i(\varphi_+ - \alpha/2)}, \quad ' \psi_- = r_- e^{i(\varphi_- - \alpha/2)}.$$

From this it follows that the two terms  $\psi_+^* \psi_-$  and  $\psi_-^* \psi_+$  transform as follows

$$'(\psi_+^* \psi_-) = r_+ r_- e^{i(\varphi - \alpha)}, \quad '(\psi_-^* \psi_+) = r_+ r_- e^{-i(\varphi - \alpha)},$$

and by (5 abc) the components of  $\Sigma$  transform as follows

$$\begin{aligned} ' \Sigma_x &= r_+ r_- (e^{-i(\varphi - \alpha)} + e^{i(\varphi - \alpha)}) \\ ' \Sigma_y &= r_+ r_- (e^{-i(\varphi - \alpha)} - e^{i(\varphi - \alpha)}) \\ ' \Sigma_z &= r_+^2 - r_-^2, \end{aligned}$$

and from these equations we conclude that

$$' \Sigma_x = 2r_+ r_- \cos(\varphi - \alpha) \quad (10a)$$

$$' \Sigma_y = 2r_+ r_- \sin(\varphi - \alpha) \quad (10b)$$

$$' \Sigma_z = r_+^2 - r_-^2. \quad (10c)$$

With this we have shown that when the wave function of the electron is transformed by a rotation with an angle  $\alpha$ , such that the two-dimensional

complex vector  $\Psi$  of its spin-components are transformed by means of the angle  $\alpha/2$  as in (6), then the vector  $\Sigma$  is rotated clockwise around the  $z$ -axis by the angle  $\alpha$  to the position given by eqs. (10 abc).

By the tandem hypothesis announced in [1], information is stored in special *storage neurons*. By definition, each storage neuron contains one and just one magnetite chain together with a prion which is switched from the *ground state* to the *excited state*, and thereby stores one bit of information when it is subject to a sufficiently strong electric impulse.

What is of interest in our present context, is the moment-vector of the magnetite chain in a storage neuron. This chain is assumed to be located close to the cell membrane in the direction of the  $y$ -axis of the chosen coordinate system, which was referred to as the “lateral direction” in [1; Sec. 6]. But in our present context it is better to use another coordinate system which is better adapted to the transformation (6) in Pauli’s theory. Therefore we will now change our coordinate system by a cyclic left-shift by which the old axis (in the forward direction of the nerve signals) is the new  $z$ -axis, while the old  $y$ -axis is the new  $x$ -axis, and the old  $x$ -axis is the new  $y$ -axis. With this change of coordinate system we get a picture different from that in Fig.1 of [1; Sec. 3], as the magnetite chain is in the direction of the  $x$ -axis, not the  $y$ -axis.

In [1; Sec. 3-4] it is explained how electrical signals in a storage neuron can trigger a boost of  $\text{Na}^+$  and  $\text{K}^+$ -ions from ion channels, which produce vortices that are rotating like blasts of a whirlwind around the magnetite chain of the neuron. These vortices initiate the quantum mechanical process which is magnetizing the chain. In this process, these positively charged ions are in rotational motion with rotation axis in the direction of the  $x$ -axis of the new coordinate system. We will assume that they are setting up a magnetic field in this direction and that the associated  $\Sigma$ -vector is ruled by equations like those in (10 abc).

More specifically, we assume that this field is the sum of contributions from the crystals in the chain, and that each of those is given (as in (2)) by a  $\Sigma$ -vector with  $x$ -component of the form

$$\Sigma_x = R \cos(\varphi - \alpha) , \tag{11}$$

with amplitude  $R > 0$  and phase angle  $\cos(\varphi - \alpha)$  where  $\varphi \in [0, 2\pi]$  and  $\alpha = \omega t$  where the rotational velocity  $\omega$  is the same for all crystals in the chain.

We will now consider a magnetite chain with  $n$  crystals, and we will denote their phase differences by  $\varphi_1, \dots, \varphi_n$  and their amplitude factors by

$R_1, \dots, R_n$ . With this, the magnetite chain in a storage neuron can be magnetized to a linear magnet in the  $x$ -direction of the chosen coordinate system, whose magnetic moment is given by the following  $x$ -component of its complex  $\Sigma$ -vector

$$\Sigma_x = \sum_{i=1}^n R_i \cos(\varphi_i - \alpha) . \quad (12)$$

When a storage neuron is in the resting state, then the crystals in the magnetite chain are completely unrelated in phase, so the phase differences  $\varphi_1, \dots, \varphi_n$  are randomly distributed and the two partial sums of positive and negative terms in (12) are likely to cancel each other out.

But when the storage neuron is subject to an electric impulse which magnetizes the chain as explained in [1; Sec. 3], then the crystals will no longer be unrelated in phase, as they are subject to a strong external field which is superimposed on the weak magnetic fields of the crystals. In this process the crystals will be forced to oscillate with increased amplitude in the same phase as the wave function of the external field, which is a quantal entity with its own amplitude and phase difference between its two spin-components.

This explains how the entire magnetite chain of a storage neuron may be magnetized by an electric impulse up to a strength which for a chain with sufficiently many crystals is large enough to switch the prion from its ground state to its excited state and thereby store one bit of information in the cell.

For especially interested readers, we will explain how equations in this note can be obtained from the general theory of quantum systems defined in abstract spaces of higher dimension than the 4-dimensional Minkowski-space.

In the formalism of quantum mechanics, a quantum system is represented on a complex Hilbert space, and the probability density of measurements of a quantum mechanical observable is represented by a positive operator of norm one, which is given by a positive matrix of trace one. If the range space of this representing operator is a single complex line, then it is the orthogonal projection onto this line. Observables for which this is the case, are said to be *in pure state*. A key example is the projection onto the line of the two-dimensional complex vector  $\Psi$  of the electron.

This complex vector  $\Psi$  is said to be a *spinor*, or more specifically a “two-dimensional spinor” since this is a special case of the general spinor concept. From general spinors we can construct, not only the rank one

matrices which determine observables in pure state, but also matrices of general observables.

The matrix of each general observable is a convex combination of matrices of observables in pure state. The matrices of observables is a convex set, and the extreme points of this convex set are the matrices of observables in pure state. For simplicity, we will identify the observables with their representing matrices. With this we will use the single word “state” to denote the matrix of a general observable and the term “pure state” to denote a the matrix of an observable in pure state. We will also use the term *state space* to denote the convex set of all matrices which represent states of the given quantum system.

The state space of a quantum system is an important concept because it determines the probability density of measurements of the observables, which is given by a “quadratic form” on the representing matrices, as in the special case of the electron, for which the position is an observable in pure state with probability density given by the matrix equation (7).

Each pair of extreme points in the state space are separated by an *angular distance* introduced in Definition 4.9 in the book [4]. These two points are pure states whose matrices are projections onto lines determined by vectors in the Hilbert space, and it can be shown that the angle between these lines (defined by the inner product of the Hilbert space) is precisely twice the angular distance between the two extreme points [4: Prop.4.14].

In this book, this result is presented in a more general context, but the proof is a simple geometric argument. What is shown, is that when a vector in the Hilbert space is rotated by the angle  $\alpha$  around the origin of the Hilbert space, then the corresponding extreme points of the state space are rotating by the double angle  $2\alpha$  around the state represented by the unit matrix. With this it is shown that the occurrence of the half-angles in equation (6) is a special case of a theorem valid for general quantum systems.

## References

- [1] Erik M. Alfsen, Fredrik C. Størmer, Arild Njå, Lars Walløe. A proposed tandem mechanism for the storage of memory in neurons.
- [2] W. Pauli. Zur Quantenmechanik des magnetischen Electrons. *Z. Phys.* nr. 43 (1927), 601–623.
- [3] Iwo Białynicki-Birula, Marek Cieplak, Jerzy Kaminsky. *Theory of Quanta*. Oxford University Press. 1992.

- [4] Erik M. Alfsen, Frederic W. Shultz. State Spaces of Operator Algebras. Birkhäuser 2001.

# A note on magnetization of magnetite chains in neurons

Erik M. Alfsen

June 11, 2018

In article [1], it is proposed that information can be stored in a neuron when a magnetized chain of magnetite crystals is reshaped and amplified by electric impulses that are switching a prion from its ground state to its excited state. In this article, this process is explained without use of mathematics. But in the present note, we will present the quantum mechanical equations for the process in which such a chain of crystals is magnetized for this purpose.

We know that quantal quantities are denoted by complex-valued functions  $\psi = re^{i\varphi}$  defined on 4-dimensional space-time (or more general real spaces not needed here), and that the absolute value  $|\psi|$  gives the probability density of the result of measurements of such quantities. We also know that quantum mechanical superposition of two quantal quantities is represented by the sum of their complex wave functions. Actually, each quantum mechanical wave function  $\psi$  consists of two spin-components, here denoted by  $\psi_+$  and  $\psi_-$ . The mathematical theory of spin dates back to Pauli's 1927-article: On the quantum mechanics of the magnetic electron [2]. Here we will give a brief survey, based on the presentation in the book [3], of those parts of Pauli's theory on electron spin which will be needed in the present note.

From the two spin-components  $\psi_+$  and  $\psi_-$  we can construct a two-dimensional vector  $\Psi$  (not in physical space but in two-dimensional complex space) and the vector conjugate to it

$$\Psi = \begin{pmatrix} \psi_+ \\ \psi_- \end{pmatrix}, \quad \Psi^\dagger = (\psi_+^*, \psi_-^*). \quad (1)$$

As explained in [2; Sec. 10.1], the magnetic moment-vector of Pauli's magnetic electron is given by an equation of the form

$$\boldsymbol{\mu} = \mu \boldsymbol{\Sigma}, \quad (2)$$



where  $\mu$  is a constant (the scalar magnetic moment), and  $\mathbf{\Sigma}$  is a three-dimensional vector (in physical space) given by the following equation for the coordinates

$$\Sigma_i = \mathbf{\Psi}^\dagger \sigma_i \mathbf{\Psi}, \quad i = 1, 2, 3 \quad \text{or} \quad i = x, y, z, \quad (3)$$

where  $\sigma_i$  are the *Pauli matrices*

$$\sigma_1 = \sigma_x = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \sigma_2 = \sigma_y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \quad \sigma_3 = \sigma_z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}. \quad (4)$$

(Here (1) is eq. (10.16) and (3) is eq. (10.12) in [3].)

Multiplying out the matrix product in (3) gives the following equations (which are (10.6a), (10.6b), (10.6c) in [3])

$$\Sigma_x = \psi_+^* \psi_- + \psi_-^* \psi_+ \quad (5a)$$

$$\Sigma_y = -i(\psi_+^* \psi_- - \psi_-^* \psi_+) \quad (5b)$$

$$\Sigma_z = \psi_+^* \psi_+ - \psi_-^* \psi_- . \quad (5c)$$

The spin-components of Pauli's magnetic electron are supposed to transform as follows under rotation by angle  $\alpha$  around a straight line in space

$${}'\psi_+ = e^{-i\alpha/2} \psi_+, \quad {}'\psi_- = e^{i\alpha/2} \psi_- . \quad (6)$$

By these equations, the wave function changes sign after rotation by the angle  $\alpha = 2\pi$ . But this won't cause problems, as it is not the wave function itself that has physical meaning, but quantities constructed from it, which are invariant under a sign-change.

The key example is the squared norm of the associated two-dimensional vector

$$|\mathbf{\Psi}|^2 = (\psi_+^*, \psi_-^*) \begin{pmatrix} \psi_+ \\ \psi_- \end{pmatrix} = |\psi_+|^2 + |\psi_-|^2, \quad (7)$$

which is interpreted as the probability density of finding the electron with given spin-state at a given point in space-time [3; eq. (10.5)]. What we are here interested in, is the transformation of the magnetic moment-vector of an electron whose wave-function is transformed as in (6).

We know that the components of the vector  $\mathbf{\Sigma}$ , which determines the magnetic moment-vector of the electron, is given as in (5). From this, we first observe that the  $z$ -component of  $\mathbf{\Sigma}$  is transformed as follows

$${}'\Sigma_z = \psi_+^* \psi_+ - \psi_-^* \psi_- = \Sigma_z . \quad (8)$$

Thus, the  $z$ -component of  $\Sigma$  remains unchanged under this transformation of the wave function. This fact is of course no mere coincidence: the transformation formula (6) was tailored to achieve just this!

To determine the transformation of  $\Sigma_x$  and  $\Sigma_y$ , we assume that the wave function of the electron at a given point  $(x, y, z)$  at a given time, say  $t = 0$ , is given by the equation

$$\psi_+ = r_+ e^{i\varphi_+}, \quad \psi_- = r_- e^{i\varphi_-}, \quad (9)$$

with the *amplitudes*  $r_+$ ,  $r_-$  and the *phase angles*  $\varphi_+$  and  $\varphi_-$ .

Note that the two terms  $\psi_+^* \psi_-$  and  $\psi_-^* \psi_+$  occurring in (5a) and (5b) are given as follows in terms of the *phase difference*  $\varphi = \varphi_+ - \varphi_-$

$$\psi_+^* \psi_- = r_+ r_- e^{-i\varphi}, \quad \psi_-^* \psi_+ = r_+ r_- e^{i\varphi},$$

and that the two terms  $\psi_+^* \psi_+$  and  $\psi_-^* \psi_-$  occurring in (5c) are given by

$$\psi_+^* \psi_+ = r_+^2, \quad \psi_-^* \psi_- = r_-^2.$$

By (6) the spin-components of an electron transform as follows under rotation by the angle  $\alpha$

$$' \psi_+ = r_+ e^{i(\varphi_+ - \alpha/2)}, \quad ' \psi_- = r_- e^{i(\varphi_- - \alpha/2)}.$$

From this it follows that the two terms  $\psi_+^* \psi_-$  and  $\psi_-^* \psi_+$  transform as follows

$$'(\psi_+^* \psi_-) = r_+ r_- e^{i(\varphi - \alpha)}, \quad '( \psi_-^* \psi_+) = r_+ r_- e^{-i(\varphi - \alpha)},$$

and by (5 abc) the components of  $\Sigma$  transform as follows

$$\begin{aligned} ' \Sigma_x &= r_+ r_- (e^{-i(\varphi - \alpha)} + e^{i(\varphi - \alpha)}) \\ ' \Sigma_y &= r_+ r_- (e^{-i(\varphi - \alpha)} - e^{i(\varphi - \alpha)}) \\ ' \Sigma_z &= r_+^2 - r_-^2, \end{aligned}$$

and from these equations we conclude that

$$' \Sigma_x = 2r_+ r_- \cos(\varphi - \alpha) \quad (10a)$$

$$' \Sigma_y = 2r_+ r_- \sin(\varphi - \alpha) \quad (10b)$$

$$' \Sigma_z = r_+^2 - r_-^2. \quad (10c)$$

With this we have shown that when the wave function of the electron is transformed by a rotation with an angle  $\alpha$ , such that the two-dimensional

complex vector  $\Psi$  of its spin-components are transformed by means of the angle  $\alpha/2$  as in (6), then the vector  $\Sigma$  is rotated clockwise around the  $z$ -axis by the angle  $\alpha$  to the position given by eqs. (10 abc).

By the tandem hypothesis introduced in [1] (and anticipated in [5]), the information is stored in special *storage neurons*. By definition, each storage neuron contains one and just one magnetite chain together with a prion which is switched from the *ground state* to the *excited state*, and thereby stores one bit of information when it is subject to a sufficiently strong electric impulse.

What is of interest in our present context, is the moment-vector of the magnetite chain in a storage neuron. This chain is assumed to be located close to the cell membrane in the direction of the  $y$ -axis of the chosen coordinate system, which was referred to as the “lateral direction” in [1; Sec. 6]. But in our present context it is better to use another coordinate system which is better adapted to the transformation (6) in Pauli’s theory. Therefore we will now change our coordinate system by a cyclic left-shift by which the old  $x$ -axis (in the forward direction of the nerve signals) is the new  $z$ -axis, while the old  $y$ -axis is the new  $x$ -axis, and the old  $z$ -axis is the new  $y$ -axis. With this change of coordinate system we get a picture different from that in Fig.1 of [1; Sec. 3], as the magnetite chain is in the direction of the  $x$ -axis, not the  $y$ -axis. (Fig.1 of this note).

In [1; Sec. 3-4] it is explained how electrical signals in a storage neuron can trigger a boost of  $\text{Na}^+$  and  $\text{K}^+$ -ions from ion channels, which produce vortices that are rotating like blasts of a whirlwind around the magnetite chain of the neuron. These vortices initiate the quantum mechanical process which is magnetizing the chain. In this process, these positively charged ions are in rotational motion with rotation axis in the direction of the  $x$ -axis of the new coordinate system. We will assume that they are setting up a magnetic field in this direction and that the associated  $\Sigma$ -vector is ruled by equations like those in (10 abc).

More specifically, we assume that this field is the sum of contributions from the crystals in the chain, and that each of those is given (as in (2)) by a  $\Sigma$ -vector with  $x$ -component of the form

$$\Sigma_x = R \cos(\varphi - \alpha) , \tag{11}$$

with amplitude  $R > 0$  and phase angle  $\varphi - \alpha$  where  $\varphi \in [0, 2\pi]$  and  $\alpha = \omega t$  where the rotational velocity  $\omega$  is the same for all crystals in the chain.

We will now consider a magnetite chain with  $n$  crystals, and we will denote their phase differences by  $\varphi_1, \dots, \varphi_n$  and their amplitude factors by

$R_1, \dots, R_n$ . With this, the magnetite chain in a storage neuron can be magnetized to a linear magnet in the  $x$ -direction of the chosen coordinate system, whose magnetic moment is given by the following  $x$ -component of its complex  $\Sigma$ -vector

$$\Sigma_x = \sum_{i=1}^n R_i \cos(\varphi_i - \alpha) . \quad (12)$$

When a storage neuron is in the resting state, then the crystals in the magnetite chain are completely unrelated in phase, so the phase differences  $\varphi_1, \dots, \varphi_n$  are randomly distributed and the two partial sums of positive and negative terms in (12) are likely to cancel each other out.

But when the storage neuron is subject to an electric impulse which magnetizes the chain as explained in [1; Sec. 3], then the crystals will no longer be unrelated in phase, as they are subject to a strong external field which is superimposed on the weak magnetic fields of the crystals. In this process the crystals will be forced to oscillate with increased amplitude in the same phase as the wave function of the external field, which is a quantal entity with its own amplitude and phase difference between its two spin-components.

This explains how the entire magnetite chain of a storage neuron may be magnetized by an electric impulse up to a strength which for a chain with sufficiently many crystals is large enough to switch the prion from its ground state to its excited state and thereby store one bit of information in the cell.

For especially interested readers, we will explain how equations in this note can be obtained from the general theory of quantum systems defined in abstract spaces of higher dimension than the 4-dimensional Minkowski-space.

In the formalism of quantum mechanics, a quantum system is represented on a complex Hilbert space, and the probability density of measurements of a quantum mechanical observable is represented by a positive operator of norm one, which is given by a positive matrix of trace one. If the range space of this representing operator is a single complex line, then it is the orthogonal projection onto this line. Observables for which this is the case, are said to be *in pure state*. A key example is the projection onto the line of the two-dimensional complex vector  $\Psi$  of the electron.

This complex vector  $\Psi$  is said to be a *spinor*, or more specifically a “two-dimensional spinor” since this is a special case of the general spinor concept. From general spinors we can construct the matrices of general observables, not only those on 4-dimensional Minkowski-space.

The matrix of each general observable is a convex combination of matrices of observables in pure state. The matrices of observables is a convex set, and the extreme points of this convex set are the matrices of observables in pure state. For simplicity, we will identify the observables with their representing matrices. With this we will use the single word “state” to denote the matrix of a general observable and the term “pure state” to denote a the matrix of an observable in pure state. We will also use the term *state space* to denote the convex set of all matrices which represent states of the given quantum system.

The state space of a quantum system is an important concept because it determines the probability density of measurements of the observables, which is given by a “quadratic form” on the representing matrices, as in the special case of the electron, for which the position is an observable in pure state with probability density given by the matrix equation (7).

Each pair of extreme points in the state space are separated by an *angular distance* introduced in Definition 4.9 in the book [4]. These two points are pure states whose matrices are projections onto lines determined by vectors in the Hilbert space, and it can be shown that the angle between these lines (defined by the inner product of the Hilbert space) is precisely twice the angular distance between the two extreme points.

In this book, this result is presented in a more general context, but the proof is a simple geometric argument: By [4; Cor.4.8] the face generated by two distinct extreme points of the state space is a Euclidean 3-ball (say of radius  $r$ ). By [4; Prop.4.11] and [4; Def.4.12] the angular distance of the two extreme points is the angle between them when viewed from the center of the ball. But the angle between them is smaller when viewed from the origin of the Hilbert space, as this point is at a longer distance ( $\sqrt{2} r$ ) from the two extreme points.

From this it follows that the angular distance between the two given extreme points (pure states) is precisely twice the angle between their representing vectors in the Hilbert space [4; Prop.14]. With this it is shown that the occurrence of the half-angles in equation (6) is a special case of a theorem valid for general quantum systems.

## References

- [1] Erik M. Alfsen, Fredrik C. Størmer, Arild Njå, Lars Walløe. A proposed tandem mechanism for the storage of memory in neurons. XXX.

- [2] W. Pauli. Zur Quantenmechanik des magnetischen Electrons. *Z. Phys.* nr. 43 (1927), 601–623.
- [3] Iwo Bialynicki-Birula, Marek Cieplak, Jerzy Kaminsky. *Theory of Quanta*. Oxford University Press. 1992.
- [4] Erik M. Alfsen, Frederic W. Shultz. *State Spaces of Operator Algebras*. Birkhäuser 2001.
- [5] Fredrik C. Størmer, Erik M. Alfsen, Is a combination of magnetite and prions involved in the memory storage in the human brain? *Medical Hypotheses* 85 (2015) 111.

